# EARLY STAGE IN THE NUCLEATION PROCESS OF CARBON NANOTUBES: DENSITY-FUNCTIONAL TIGHT-BINDING MOLECULAR DYNAMICS SIMULATIONS OF ACCETYLENE OLIGOMERIZATION AND CROSS-LINKING ON AN FE<sub>38</sub> PARTICAL

Ying Wang,<sup>1</sup> Yasuhito Ohta,<sup>2</sup> HuJun Qian,<sup>1</sup> Keiji Morokuma<sup>3</sup> and Stephan Irle<sup>1</sup>

<sup>1</sup>Institute for Advanced Research and Department of Chemistry, Nagoya University, Nagoya 464-8602, Japan <sup>2</sup>Department of Chemistry Faculty of Science, Nara Women University, Nara 630-8506, Japan <sup>3</sup>Fukui Institute for Fundamental Chemistry, Kyoto University, Kyoto, 606-8103, Japan

#### Introduction

Catalytic chemical vapor deposition (CCVD) is typically used to synthesize carbon nanotubes (CNTs), but the exact role of metal catalysts is still unknown. Since Fe substrate is known to be highly effective for the CNT growth in the CCVD process and  $C_2H_2$  is an excellent carbon source for high yield [1-5], we performed quantum chemical molecular dynamics (QM/MD) simulations based on the density-functional tight-binding (DFTB) method [6] to investigate the self-assembly process of carbon nanotubes from acetylene molecules on an Fe<sub>38</sub> particle. We found that oligomerization and cross-linking reactions between carbon chains with high H:C ratio were observed as initial steps, preceding any future cap nucleation.

## Computational methodology A. Density-Functional Tight-Binding method (DFTB)

The density-functional tight binding (DFTB) method is the central method adopted in the present studies. All DFTB calculations were carried out with DFTB program package developed by Frauenheim, Seifert, and Elstner [6-8]. DFTB is an approximate density functional theory method based on the tight binding approach and utilizes an optimized minimal LCAO Slater-type all valence basis set in combination with two-center approximation for Hamiltonian matrix elements. The DFTB method is approximately two orders of magnitude faster than first-principles density functional theory (DFT) but often yields comparable accuracy, therefore enables longer simulations and more adequate model systems for nonequilibrium dynamics of nano-sized clusters with quantum mechanical treatment of the electronic degrees of freedom. The molecular orbitals obey the following generalized eigenvalue equation,

$$\sum_{v} c_{vi} \left( H_{\mu v} - \varepsilon_i S_{\mu v} \right) = 0$$

and the total energy is described as follows,

$$E_{tot} = 2\sum_{i} f_{i}\varepsilon_{i} + E_{rep} + \frac{1}{2}\sum_{\alpha\beta}\gamma_{\alpha\beta}\Delta q_{\alpha}\Delta q_{\beta}$$

Where  $H_{\mu\nu}$  and  $S_{\mu\nu}$  are Hamiltonian matrix and overlap matrix, respectively,  $f_i$  and  $\varepsilon_i$  represent fractional occupation and orbital energy of the *i*<sup>th</sup> Kohn-Sham (KS) molecular orbital,  $E_{rep}$  denotes the distance-dependent diatomic repulsive potentials, and  $\gamma_{\alpha\beta}$  represents the distance-dependent chargecharge interaction functional, which is derived from the chemical hardness (IP–EA). The  $\Delta q_{\alpha}$  are induced charges on atom  $\alpha$ . In the SCC-DFTB method, the atomic charge-charge interaction is explicitly included in the electronic energy, and the DFTB secular equation is solved iteratively until the Mulliken atomic charges of the KS orbitals become selfconsistent with the Hamiltonian.

## **B.** Molecular Dynamics simulation

The DFTB/MD approach is a Born-Oppenheimer molecular dynamics technique based on the DFTB electronic structure method, and can bridge the gap between classical and the first principle MD simulations. In the present DFTB/MD simulations, the self-consistent-charge DFTB (SCC-DFTB) method in combination with a finite electronic temperature approach [9-11] with Te = 10000 K was employed to evaluate the quantum chemical potential on the fly. The electronic temperature allows the occupancy of each molecular orbital to change smoothly from 2 to 0 depending on its energy, and effectively incorporates the open-shell nature of the system due to near-degeneracy of iron d orbitals as well as carbon dangling bonds. In the molecular dynamics simulations the velocity Verlet integrator was used with a short time step of 0.5 fs. The nuclear temperature (Tn=1500 K) was controlled by connecting the Nose-Hoover chain thermostat [12] to the degrees of freedom of the present model system. Periodic boundary conditions were imposed with a cubic box size of 100 Angstrom<sup>3</sup>.

#### **Results and Discussion**

In our MD simulations, we first equilibrated an fcc-Fe<sub>38</sub> icosahedral magic number cluster at 1500 K for 10 ps. Then, ten structures were randomly chosen from this equilibration run between the 5 and 10 ps mark. Using these 10 Fe<sub>38</sub> initial structures, we started to "shoot" single  $C_2H_2$ molecules every 1 ps onto the metal cluster until their total number was 30. During the beginning of these 30 ps carbon supply simulations we found that H occasionally transfers to the iron particle or other  $C_2H_2$  units on the cluster to form C<sub>2</sub>H<sub>3</sub> species, producing C<sub>2</sub>H. Subsequently, these more reactive radical species initiate oligomerization reactions among the C2-containing molecules. In the following, we annealed the resulting clusters for 80 ps at the same temperature. During these carbon diffusion simulations, five-(trajectories A, E, and I) or six-membered rings (Trajectory F), as well as longer carbon chains were formed from oligomerization or cross-linking. All trajectories produce

pentagons first. For trajectory E, and F, the ring formed from a single chain  $(C_6H_x \text{ or } C_8H_x)$ , similar to trajectory A, as shown in Fig. 1b, while for trajectory I, the ring formation mechanism results from cross-linking between C<sub>6</sub>H<sub>x</sub> or C<sub>4</sub>H<sub>y</sub>, as shown in Fig. 2. This cross-linking is similar to the mechanism proposed in Reference [1], however, we did not observe the direct formation of H<sub>2</sub>, and cross-linking did not form a six-membered ring network directly. We show here the analysis of trajectory A as a representative example, and give the carbon hybridization state analysis and cluster size histogram in Figs. 1 c-d. We found that the numbers of double- (sp<sup>2</sup>) and triple- (sp) bonded carbons are of similar magnitude. Fig. 1d shows that the number of  $C_2H_x$  increases gradually until about 45 carbons are present on the metal cluster. With increasing time, larger clusters (blue, red and yellow areas) gradually consume smaller fragments.

Accompanying our QM/MD simulations, we computed the barrier heights for four kinds of reactions at the DFTB level: H-transfer to Fe (~40 kcal/mol), C-C bond breaking (~45 kcal/mol), H<sub>2</sub> abstraction (~70 kcal/mol), and C-C oligomerization (~23 kcal/mol). These energetics explain why C-C oligomerization occurs readily since it is an exothermic reaction (~ -10 kcal/mol). To the contrary, direct H<sub>2</sub> abstraction possesses a very high barrier and is an endothermic process (~68 kcal/mol), less likely to occur in MD.

# Conclusions

During 30 ps Carbon supply simulation, some of  $C_2H_2$  loose H or some H transfer to C to form  $C_2H_x$  and oligomerization occurs, especially the connection between  $C_2H_2$  and  $C_2H$  (8/10 ratio), which has a lower barrier of 23.39 kcal/mol and it is an exothermic reaction (15.13 kcal/mol). During 80 ps carbon diffusion simulation, some five or six-membered rings formed from long chain  $C_6H_x$  or  $C_8H_x$ , and by cross-linking. We found pentagon-first mechanism. H<sub>2</sub> abstraction possesses a very high barrier and is an endothermic process, then it was consequently never observed in our 110 ps MD simulations. We hope our results can serve as a basis for future study of growth mechanism of CNT as well as catalytic decomposition of hydrocarbon on metal surface.

#### References

- [1] Eres G, Rouleau CM, Yoon M, Puretzky AA, Jackson JJ,
- Geohegan DB. J. Phys. Chem. 2009; 113并 1548
- [2] Satishkumar BC, Govindaraj A, Sen R, Rao CNR. Chem. Phys. Lett 1998; 293: 47.
- [3] Ci L, Xie S, Tang D, Yan X, Li Y, Liu Z, Zhou X, Zhou W, Wang C, Cham Phys. Lett 2001; 240; 101
- Wang G, Chem. Phys. Lett 2001; 349: 191.
- [4] Liu BC, Lyu SC, Jung SI, Kang HK, Yang CW, Park JW, Park CY, Lee CJ. Chem. Phys. Lett. 2004; 383: 104.
- [5] Yamaguchi M, Pan L, Akita S, Nakayama Y. J Applied Phys. 2008; 47: 1937.
- [6] Elstner M, Porezag D., Jungnickel G, Elsner J, Haugk M,
- Frauenheim T, Suhai S, Seifert G. Phys. Rev. B 1998; 58: 7260.

[7] Porezag D, Frauenheim T, Seifert G, Kaschner, R. Phys. Rev. B. 1995; 51: 12947.

- [8] Seifert G, Porezag D, Frauenheim T, Suhai S. Int. J. Quantum Chem. 1996; 58: 185.
- [9] Weinert M, Davenport JW. Phys. Rev. B 1992; 45 (23): 13709.
- [10] Wentzcovitch RM, Martins JL, Allen PB. Phys. Rev. B 1992; 45 (19): 11372.
- [11] Wagner F, Laloyaux T, Scheffler M. Phys. Rev. B 1998, 57 (4): 2102.
- [12] Martyna GJ, Klein ML, Tuckerman M. J. Chem. Phys. 1992; 97: 2635.



**Fig. 1** Analysis of trajectory A. Carbons undergoing chemical transformations are highlighted by purple and blue spheres. a) The snapshots after carbon supply simulation and carbon diffusion simulation; b) ring count statistics as a function of time and schematic depictions of pentagons formation procession from polyyne-like chains; c) carbon hybridization analysis; d) carbon cluster size count.



**Fig. 2** Crossing Link for trajectory I. Reactive carbons are highlighted by purple and blue spheres, 1 and 2 indicate first and second connection.